

Similarity of mass, momentum, and heat transport laws

(Empirical laws)

Fick's first law : diffusion of mass

$$J_{ix} = -D \frac{dC_i}{dx} \quad \begin{array}{l} J: [ML^{-2} T^{-1}] \\ D : [L^2 T^{-1}] \end{array}$$

Newton's law : dissipation of momentum (force) through fluid viscosity

$$\tau_{yx} = -\mu \frac{d u_x}{dy} \quad \tau_{yx} : [ML^{-1} T^{-2}]$$

fluid shear

$$= - \frac{\mu}{\rho} \frac{d u_x}{dy} \quad \nu = \frac{\mu}{\rho} [L^2 T^{-1}]$$

kinematic viscosity

Fourier's law : heat conduction

$$q_x = -\kappa \frac{dT}{dx} \quad q : [cal M^{-1} T^{-1}]$$

$$= - \frac{\kappa}{C_p \rho} \frac{dT}{dx} \quad \alpha = \frac{\kappa}{C_p \rho} [L^2 T^{-1}]$$

thermal diffusivity

Table 4.1

Diffusion coeff	order of magnitude		
	Air	Water	
ν_j	0.2	0.01	(momentum)
α	0.2	1.5×10^{-3}	(heat)
D_{ij}	0.2	10^{-5}	(diffusion)

why equal? why different?

$$Sc = \frac{\text{momentum diffusivity } \nu}{\text{chemical diffusivity } D} = \frac{\nu}{D} \quad \text{Schmidt number}$$

$$Pr = \frac{\text{momentum diffusivity } \nu}{\text{thermal diffusivity } \alpha} = \frac{\nu}{\alpha} \quad \text{Prandtl number}$$

for air (or low pressure gaseous systems) $Pr \approx Sc \approx 0.7$

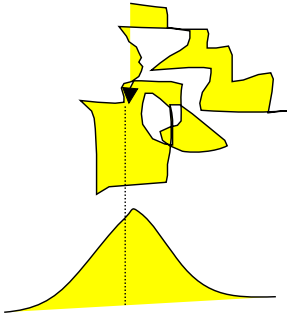
Principle of Diffusion : random molecular motion

→ complete mixing in a confined space
thermally induced fluctuation

translational kinetic energy = $\frac{3}{2} kT = \frac{1}{2} mU$

probability of a displacement

between x and $x + dx$ after n random steps of length is given by the Gaussian distribution



$$P(x, t) = \frac{1}{\sqrt{2\pi n l^2}} e^{-x^2/2nl^2} dx$$

$$n = kt$$

$$C = C_0 P(x, t)$$

This is the solution of one-dimensional diffusion eqn

$$\frac{\partial C}{\partial t} = D \frac{\partial^2 C}{\partial x^2}$$

$$C = C_0 \frac{1}{\sqrt{2\pi Dt}} e^{-x^2/4Dt}$$

in gases	diffusion progress about	10 cm min
liquid		0.05 cm min
solid		0.00001 cm min

Diffusion : random molecular motion (Brownian motion)

→ complete mixing

slowest step in the sequence

compared to convective flow, reaction, mixing(stirring)

wind 100 m/min

but determine the efficiency of a unit operation

◀ mixing ~ hydrodynamic movement
diffusion ~ molecular movement

$$\text{flux} = \frac{\text{amount transported}}{(\text{time}) (\text{area})}$$

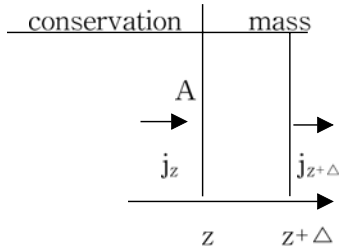
Fick's second law (conservation of mass)

Fick's first law

total flux $J = Aj = \frac{\partial}{\partial t} AD$

one-dimensional diffusion equation (empirical equation)

why $\frac{\partial}{\partial t} > 0, j_1 < 0$



$$\frac{\partial}{\partial t} \frac{A \cdot \Delta}{\text{volume}} = (j_z - j_{z+\Delta}) A$$

$$\frac{\partial}{\partial t} = - \frac{j_{z+\Delta} - j_z}{\Delta}$$

$$= - \frac{\partial}{\partial t}$$

$$\frac{\partial}{\partial t} = + D \frac{\partial^2 c}{\partial z^2} : \text{one-dimensional Fick's second law}$$

3-dimensional or vector form of Fick's law

$$j = -D \nabla c \qquad j = -D \left(\frac{\partial}{\partial x} + \frac{\partial}{\partial y} + \frac{\partial}{\partial z} \right)$$

$$\frac{\partial}{\partial t} = + D \nabla^2 c \qquad \frac{\partial}{\partial t} = D \left(\frac{\partial^2 c}{\partial x^2} + \frac{\partial^2 c}{\partial y^2} + \frac{\partial^2 c}{\partial z^2} \right)$$

case 1. A is not constant $\frac{\partial}{\partial t} = D \left(\frac{\partial^2 c}{\partial z^2} + \frac{1}{A} \frac{\partial A}{\partial z} \frac{\partial c}{\partial z} \right)$

case 2. D is directional $\frac{\partial}{\partial t} = D_x \frac{\partial^2 c}{\partial x^2} + D_y \frac{\partial^2 c}{\partial y^2} + D_z \frac{\partial^2 c}{\partial z^2}$

general (constitutive) transport equation

for constant D

$$\frac{\partial}{\partial t} = -(\mathbf{u} \cdot \nabla c - \nabla \cdot (-D \nabla c) + R_c) \Rightarrow \frac{\partial}{\partial t} + \mathbf{u} \cdot \nabla c - D \nabla^2 c - R_c = 0$$

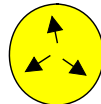
coordinates of Fick's law

$$j_1 = -D \frac{dq}{dx} \quad \text{one-dimensional cartesian coordinate}$$

$$j_r = -D \frac{dq_r}{dr} \quad \text{radial direction in cylindrical coordinate}$$



$$j_r = -D \frac{dq_r}{dr} \quad \text{radial direction in spherical coordinate}$$

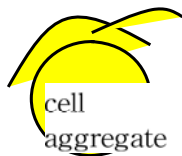


2nd law is different for each coordinate

$$\frac{\partial}{\partial t} = D \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial}{\partial r} \right) \quad \text{cylindrical}$$

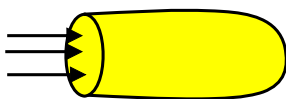
$$\frac{\partial}{\partial t} = D \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) \quad \text{spherical}$$

* coordinate is not determined by the geometry itself.



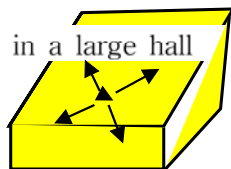
analysis of
thin layer

→ can be simplified to one-dimensional cartesian coordinate



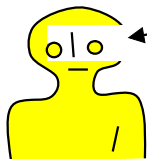
transport
in a capillary
tube

→ " "



in a large hall

diffusion from a point source
in a large → cartesian or spherical
room coordinate



acid
rain

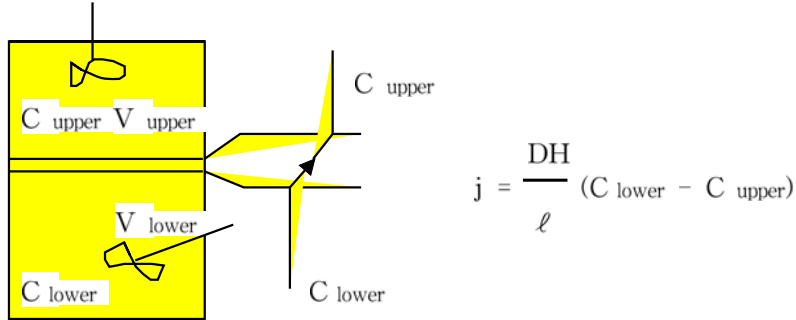
very slow
diffusion

→ one dimensional cartesian

Diffusion Coefficient

Air-CO ₂	0.142 cm ² sec	at 276.2 K
Air-benzene	0.096 cm ² sec	at 298 K
O ₂ in water	2.10 × 10 ⁻⁵ cm ² sec	at 298 K
Cl ₂ in water	1.25 × 10 ⁻⁵ cm ² sec	at 298 K

Experimental Determination of Diffusion Coefficient



rate of decrease in the lower compartment

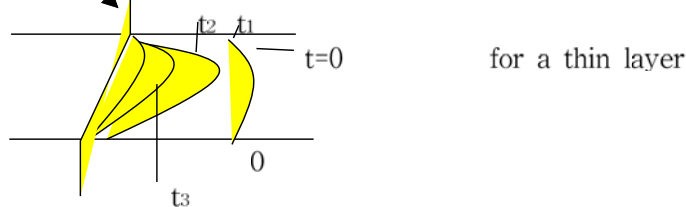
$$V_{\text{lower}} \frac{dC_{\text{lower}}}{dt} = -Aj, \quad \frac{dC_{\text{lower}}}{dt} = -\frac{A}{V_{\text{lower}}} j \quad (1)$$

rate of increase in the upper compartment

$$V_{\text{upper}} \frac{dC_{\text{upper}}}{dt} = Aj, \quad \frac{dC_{\text{upper}}}{dt} = \frac{A}{V_{\text{upper}}} j \quad (2)$$

Assume quasi steady state →

$$j = \frac{DH}{l} (C_{\text{lower}} - C_{\text{upper}})$$



$$\frac{d}{dt} (C_{\text{lower}} - C_{\text{upper}}) = -\left(\frac{A}{V_{\text{lower}}} + \frac{A}{V_{\text{upper}}}\right) \frac{DH}{l} (C_{\text{lower}} - C_{\text{upper}})$$

$$y \equiv C_{\text{lower}} - C_{\text{upper}} \quad \beta \left(\frac{A}{V_{\text{lower}}} + \frac{A}{V_{\text{upper}}}\right) \frac{DH}{l}$$

initial concentration

$$\frac{dy}{dt} = -\beta y \quad \text{at } t=0 \quad y = C_{\text{lower}} - C_{\text{upper}} = y^0$$

$$y = y^0 e^{-\beta t}$$

$$\frac{C_{\text{lower}} - C_{\text{upper}}}{C_{\text{lower}} - C_{\text{upper}}} = e^{-\beta}$$

experimental data (unsteady state)
 → data regression
 → $\beta \rightarrow \infty$

H : distribution coefficient
 H ~ 1 preferred (chemically inert membrane-separator)

General transport equations for various coordinates

Table 3.4-2. Mass balance for species 1 combined with Fick's law

Rectangular coordinates	
$\frac{\partial c_1}{\partial t} + v_x^0 \frac{\partial c_1}{\partial x} + v_y^0 \frac{\partial c_1}{\partial y} + v_z^0 \frac{\partial c_1}{\partial z} = D \left(\frac{\partial^2 c_1}{\partial x^2} + \frac{\partial^2 c_1}{\partial y^2} + \frac{\partial^2 c_1}{\partial z^2} \right) + r_1$	(A)
Cylindrical coordinates	
$\frac{\partial c_1}{\partial t} + v_r^0 \frac{\partial c_1}{\partial r} + \frac{v_\theta^0}{r} \frac{\partial c_1}{\partial \theta} + v_z^0 \frac{\partial c_1}{\partial z} = D \left[\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial c_1}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 c_1}{\partial \theta^2} + \frac{\partial^2 c_1}{\partial z^2} \right] + r_1$	(B)
Spherical coordinates	
$\frac{\partial c_1}{\partial t} + v_r^0 \frac{\partial c_1}{\partial r} + v_\theta^0 \frac{\partial c_1}{\partial \theta} + \frac{v_\phi^0}{r \sin \theta} \frac{\partial c_1}{\partial \phi}$	
$= D \left[\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial c_1}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial c_1}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 c_1}{\partial \phi^2} \right] + r_1$	(C)

Note: The diffusion coefficient D and the density ρ are assumed constant. In this case, the mass average and volume average velocities are equal. Again, r_1 is the rate of production of species 1 per volume.

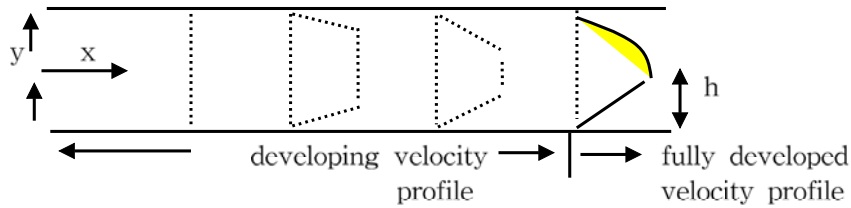
Table 3.4-3. Total mass balance in several coordinate systems

Rectangular coordinates	
$\frac{\partial \rho}{\partial t} = -\frac{\partial}{\partial x}(\rho v_x) - \frac{\partial}{\partial y}(\rho v_y) - \frac{\partial}{\partial z}(\rho v_z)$	(A)
Cylindrical coordinates	
$\frac{\partial \rho}{\partial t} = -\frac{1}{r} \frac{\partial}{\partial r}(\rho r v_r) - \frac{1}{r} \frac{\partial}{\partial \theta}(\rho v_\theta) - \frac{\partial}{\partial z}(\rho v_z)$	(B)
Spherical coordinates	
$\frac{\partial \rho}{\partial t} = -\frac{1}{r^2} \frac{\partial}{\partial r}(\rho r^2 v_r) - \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta}(\rho v_\theta \sin \theta) - \frac{1}{r \sin \theta} \frac{\partial}{\partial \phi}(\rho v_\phi)$	(C)

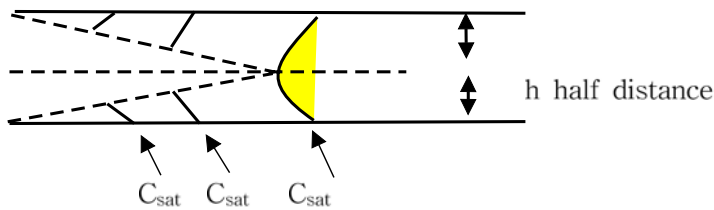
Note: The velocity here is the mass average and not the volume average commonly used with Fick's law.

General Transport Equations and Their Applications
Two Case Studies on Analytical Solutions of Partial Differential Equations in Mass Transfer using Similarity Variables

Case 1 Mass transfer in a channel flow
 flow development in a channel



development of concentration profile in a channel
 with slightly soluble walls (no change in gap)



" Concentration profile in a fully developed flow"
 velocity profile for a fully developed flow

Navier Stokes eqn

$$\rho \frac{Du}{Dt} = \mu \nabla^2 u - \nabla p + \rho g \quad (4-17)$$

steady state $Du/Dt = 0$

No gravity $\rho g = 0$

Δp is the only driving force

$$\nabla^2 u = \frac{\Delta p}{\mu} \quad (1)$$

$$\frac{d^2 u}{dy^2} = \frac{\Delta p}{\mu} \quad \text{BC I } y=0 \quad \frac{du}{dy}=0 \quad (2)$$

$$\text{BC II } y=h \quad u=0$$

$$\frac{du}{dy} = \frac{\Delta p}{\mu} y + C_1 \quad \text{BC I. } 0=C_1 \quad (3)$$

$$u = \frac{\Delta p}{2\mu} y^2 + C_1 y + C_2 \quad \text{BC II} \quad 0 = \frac{\Delta p}{2\mu} h^2 + C_2 \quad (4)$$

$$u = -\frac{\Delta}{2\mu}(h^2 - y^2) \quad C_2 = \frac{\Delta}{2\mu} h^2 \quad (5)$$

$$u = -\frac{\Delta}{2\mu} \left[1 - \left(\frac{y}{h}\right)^2 \right] \quad r=0 \quad u_{\max} = \frac{\Delta}{\mu} \quad (6)$$

near the wall

$$1 - (y/h)^2 \quad 1 - [(y^2 - 2yh + h^2)/h^2] \approx 2y/h \quad y^2 \ll h^2 \quad (7)$$

$$u = u_{\max}(2y/h) \quad (8)$$

(how to derive from force balance ?)

Constitutive transport equation

$$\left(\frac{\partial}{\partial t} + \mathbf{u} \cdot \nabla \right) C_i = D \nabla^2 C_i + R_i \quad (4-28)$$

steady state $\frac{\partial}{\partial t} = 0$

no reaction $R_i = 0$

$$\mathbf{u} \cdot \nabla = D \nabla^2 C \quad (9)$$

for cartesian coordinate

$$u_x \frac{\partial}{\partial x} + u_y \frac{\partial}{\partial y} + u_z \frac{\partial}{\partial z} = D \frac{\partial^2 C}{\partial x^2} + D \frac{\partial^2 C}{\partial y^2} + D \frac{\partial^2 C}{\partial z^2}$$

$$\frac{2y}{h} u_{\max} \frac{\partial}{\partial x} = D \frac{\partial^2 C}{\partial x^2} \quad \text{two independent var. } x, y \quad \text{one dependent var, } c \quad (10)$$

$$\begin{aligned} C &= C_{\text{sat}} & \text{at } y = 0 \\ C &= 0 & \text{as } y \rightarrow \infty \quad ** \quad \frac{\partial C}{\partial y} = 0 \quad \text{at } y=h \end{aligned} \quad (11)$$

$$\frac{C}{C_{\text{sat}}} = C^* \quad f(\eta) \quad \left(\eta = y, \frac{u_{\max}}{Dh} \right)$$

Similarity : non-dimensional combination
similarity variable

$$\eta = y \left(\frac{2u_{\max}}{hDx} \right)^{1/3} \quad (12)$$

chain rule $C^* = f(\eta)$

$$\frac{\partial C^*}{\partial x} = \frac{dC^*}{d\eta} \frac{\partial \eta}{\partial x} \Big|_{y \text{ const}} = y \left(\frac{2u_{\max}}{hD} \right)^{1/3} \left(-\frac{1}{3} \right) x^{-4/3} \frac{dC^*}{d\eta} = -\frac{\eta}{3x} \frac{dC^*}{d\eta} \quad (13)$$

$$\frac{\partial C^*}{\partial y} = \frac{dC^*}{d\eta} \frac{\partial \eta}{\partial y} \Big|_{x \text{ const}} = \left(\frac{2u_{\max}}{hDx} \right)^{1/3} \frac{dC^*}{d\eta} \quad (14)$$

$$\frac{\partial^2 C^*}{\partial x^2} = \left(\frac{2u_{\max}}{hDx} \right)^{1/3} \frac{dC^*}{d\eta} \frac{\partial \eta}{\partial x} = \left(\frac{2u_{\max}}{hDx} \right)^{2/3} \left(\frac{d^2 C^*}{d\eta^2} \right) \quad (15)$$

$$\cancel{\frac{y}{h}} u_{\max} \left(\frac{\eta}{3x} \right) \frac{dC^*}{d\eta} = D \left(\frac{2u_{\max}}{hDx} \right)^{2/3} \frac{dC^*}{d\eta} \quad (16)$$

$$\frac{dC^*}{d\eta} + \frac{\eta}{3} \frac{dC^*}{d\eta} = 0 \quad (17)$$

BC $C^* = 1$ at $\eta = 0 \leftarrow y = 0$
 $C^* = 0$ at $\eta = \infty \leftarrow y \rightarrow \infty$ (18)

$$\frac{d}{d\eta} \ln \left(\frac{dC^*}{d\eta} \right) = -\frac{\eta}{3} \quad (19)$$

$$C^* = A \int_0^{\eta} e^{-\eta^2/9} d\eta + B \quad (20)$$

$$C^* = 1 - \frac{\int_0^{\eta} e^{-\eta^2/9} d\eta}{\int_0^{\infty} e^{-\eta^2/9} d\eta} \quad (21)$$

Gamma function $\Gamma(\nu) = \int_0^{\infty} Z^{\nu-1} e^{-Z} dZ \quad (\nu > 0)$ (22)

$$\Gamma_{i+1} = \int_0^\infty Z e^{-z^2} dz \quad \Gamma_{i+1} = \pi \Gamma_i$$

$$\frac{C}{C_{\text{sat}}} = 1 - 0.538 \int_0^{\eta} e^{-\eta^2/9} d\eta \quad (23)$$

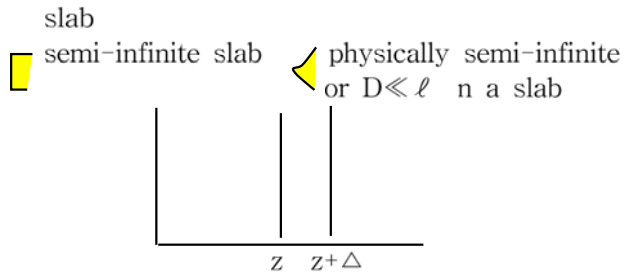
Flux at wall

$$j^* = -D \left(\frac{\partial}{\partial y} \right)_w = -C_{\text{sat}} D \left(\frac{2u_{\text{max}}}{hDx} \right)^{1/3} \left(\frac{dC^*}{d\eta} \right)_w \quad (y=h) \quad (24)$$

$$= 0.678 C_{\text{sat}} D \left(\frac{u_{\text{max}}}{hDx} \right)^{1/3} \quad (25)$$

Equation (25) is useful, since the solubility (or total flux) is expressed as a function of distance. Also it is coincident with the empirical relations for mass transfer.

Case 2 Unsteady state diffusion in a semi-infinite slab starting from mass balance (free diffusion)



$$\left(\begin{array}{c} \text{solute accumulation} \\ \text{in volume } A\Delta \end{array} \right) = \left(\begin{array}{c} \text{rate of diffusion} \\ \text{into the layer} \\ \text{at } Z \end{array} \right) - \left(\begin{array}{c} \text{rate of diffusion} \\ \text{out of the layer} \\ \text{at } Z+\Delta \end{array} \right)$$

$$\frac{\partial}{\partial t} (A\Delta C_1) = A \left(j_1 \Big|_Z - j_1 \Big|_{Z+\Delta} \right)$$

$$\frac{\partial C_1}{\partial t} = - \frac{j_1 \Big|_{Z+\Delta} - j_1 \Big|_Z}{(Z+\Delta) - Z}$$

$$\begin{array}{l} t=0 \quad \text{all } z \quad C_1 = C_1 \infty \\ t \rightarrow \quad z=0 \quad C_1 = C_1 0 \\ \quad \quad z=\infty \quad C_1 = C_1 \infty \end{array} \quad \frac{\partial C_1}{\partial t} = - \frac{\partial j_1}{\partial z} \quad j_1 = -D \frac{\partial C_1}{\partial z}$$

$$\frac{\partial C_1}{\partial t} = D \frac{\partial^2 C_1}{\partial z^2} \quad \leftarrow \text{Fick's 2nd law Diffusion equation (1)}$$

\leftarrow Boundary condition

$$\zeta \equiv \frac{z}{\sqrt{4Dt}} \quad \text{combination of variables (2)}$$

$$\frac{dC_1}{d\zeta} \left(\frac{\partial \zeta}{\partial t} \right) = D \frac{d^2 C_1}{d\zeta^2} \left(\frac{\partial \zeta}{\partial z} \right) \quad (3)$$

$$\text{LHS : } \frac{d\zeta}{dt} = \frac{d}{dt} \left(\frac{z}{\sqrt{4Dt}} \right)^{1/2} = \frac{z}{4Dt} \quad (4)$$

$$\text{RHS : } \left(\frac{d\zeta}{dz} \right)^2 = \left(\frac{1}{\sqrt{4Dt}} \right)^2 = \frac{1}{4Dt} \quad (5)$$

$$\frac{dC_1}{d\zeta} \left(-\frac{2z}{4t\sqrt{Dt}} \right) = D \frac{d^2C_1}{d\zeta^2} - \frac{1}{4Dt} \quad (6)$$

$$\frac{d^2C_1}{d\zeta^2} + \frac{2z}{\sqrt{Dt}} \frac{dC_1}{d\zeta} = 0 \quad (7)$$

$$\frac{d^2C_1}{d\zeta^2} + 2\zeta \frac{dC_1}{d\zeta} = 0 \quad (8)$$

$$\begin{aligned} z=0 \quad C_1 = C_{10} &\Rightarrow \zeta = 0 \quad C_1 = C_{10} \\ t=0 \\ z=\infty \quad C_1 = C_{1\infty} &\Rightarrow \zeta = \infty \quad C_1 = C_{1\infty} \end{aligned} \quad (9)$$

integrate

$$\frac{dC_1}{d\zeta} = a e^{-\zeta}$$

confirm

$$\frac{d}{d\zeta} C_1 = a e^{-\zeta} \quad (-2\zeta)$$

$$\text{integrate } \int dC_1 = \int a e^{-\zeta} d\zeta$$

$$C_1 - C_{10} = a \int_0^x e^{-y} dy$$

independent variable?

$$C_1 - C_{10} = a \frac{\sqrt{\pi}}{2} \operatorname{erf}\zeta \quad (11)$$

a=?

at $\zeta = \infty \quad C_1 = C_{1\infty}$

$$C_{1\infty} - C_{10} = a \frac{\sqrt{\pi}}{2} \operatorname{erf}\infty \quad \operatorname{erf}\infty = 1$$

$$= \frac{a\sqrt{\pi}}{2}$$

$$a = (C_{1\infty} - C_{10}) \frac{2}{\sqrt{\pi}}$$

$$\begin{aligned} &= a \frac{1}{\sqrt{\pi}} \frac{dC_1}{d\zeta} (-2\zeta) \\ &\frac{d}{d\zeta} C_1 + 2\zeta \frac{dC_1}{d\zeta} = 0 \quad (10) \end{aligned}$$

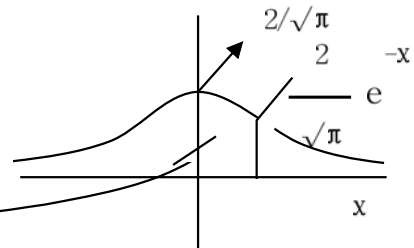
$$\operatorname{erf}x \equiv \frac{2}{\sqrt{\pi}} \int_0^x e^{-y} dy$$

$$C_1 - C_{10} = (C_{1\infty} - C_{10}) \frac{\sqrt{\pi}}{2} \frac{2}{\sqrt{\pi}} \operatorname{erf} \zeta$$

$$\frac{C_1 - C_{10}}{C_{1\infty} - C_{10}} = \operatorname{erf} \zeta \quad (12)$$

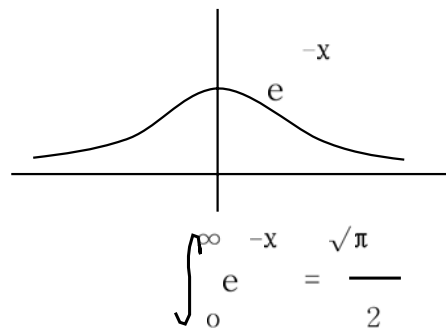
Properties of error function

$$\operatorname{erf} x = \frac{2}{\sqrt{\pi}} \int_0^x e^{-y^2} dy$$



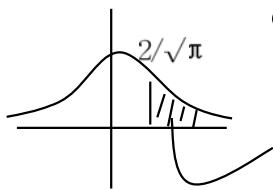
$$\begin{cases} \operatorname{erf}(\infty) = 1 \\ \operatorname{erf} 0 = 0 \end{cases}$$

$$\frac{d}{dx} \operatorname{erf} x = \frac{2}{\sqrt{\pi}} e^{-x^2}$$



$$\int_0^{\infty} e^{-x^2} dx = \frac{\sqrt{\pi}}{2}$$

$$\operatorname{erfc} x = 1 - \operatorname{erf} x$$



$$\operatorname{erfc}(x) = \frac{2}{\sqrt{\pi}} \int_x^{\infty} e^{-y^2} dy$$

$$\begin{cases} \operatorname{erfc} 0 = 1 \\ \operatorname{erfc} \infty = 0 \end{cases}$$

Flux

$$j_1 = -D \frac{\partial c_1}{\partial z} = -D(C_{1\infty} - C_{10}) \frac{2}{\sqrt{\pi}} e^{-\zeta^2} \frac{d\zeta}{dz} \quad (13)$$

$$= -D(C_{1\infty} - C_{10}) \frac{2}{\sqrt{\pi}} e^{-Z^2/4Dt} \frac{1}{\sqrt{4Dt}}$$

$$= (C_{10} - C_{1\infty}) \frac{\sqrt{D}}{\sqrt{\pi}} e^{-Z^2/4Dt} \quad (14)$$

flux at z=0

$$j_1 \Big|_{z=0} = (C_{1\infty} - C_{10}) \frac{\sqrt{D}}{\sqrt{\pi}} \quad (15)$$

Superposition method

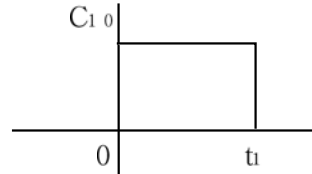
$$\frac{C_1 - C_{10}}{C_{1\infty} - C_{10}} = \text{erf}\zeta$$

$C_{1\infty} = 0$ (or initial condition)

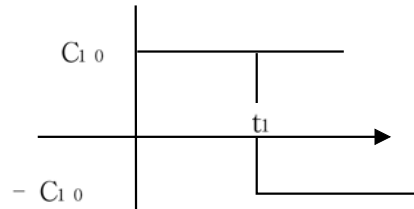
$$1 - C_1/C_{10} = \text{erf}\zeta$$

$$C_1/C_{10} = 1 - \text{erf}\zeta = \text{erfc}\zeta$$

$$\begin{cases} 0 < t < t_1 & C = C_{10} \\ t_1 \leq & C = 0 \end{cases}$$



Du Hamel's theorem

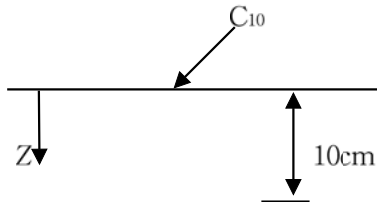


$$\begin{aligned} C_1 &= C_{10} \text{erfc}\zeta \quad (0 < t < t_1) \\ C_1 &= C_{10} \text{erfc}\zeta - C_{10} \text{erfc}\zeta_1 \end{aligned}$$

$$\zeta = \frac{z}{\sqrt{Dt}} \quad \zeta_1 = \frac{z}{\sqrt{D(t-t_1)}}$$

$$C_1 = C_{10} (1 - \text{erf}\zeta - 1 + \text{erf}\zeta_1) = C_{10} (\text{erf}\zeta_1 - \text{erf}\zeta)$$

Example (soil contamination)



semi infinite

at what time
C at 10cm = 0.02 kg/kg soil

$$\frac{\partial C_A}{\partial t} = D \frac{\partial^2 C_A}{\partial z^2}$$

$$\begin{aligned} C_{\text{sat}} &= 0.05 \text{ kg/kg soil} \\ D_{\text{eff}} &= 5 \times 10^{-4} \text{ m}^2/\text{sec} \end{aligned}$$

$$\begin{aligned} \text{IC } t=0 \text{ for all } Z & C_A = C_{A\infty} \\ t>0 \quad Z=0 & C_A = C_{A0} = 0.05 \text{ kg/kg} \\ Z=\infty & C_A = C_{A\infty} = 0 \end{aligned}$$

$$\frac{C_A - C_{A0}}{0 - C_{A0}} = \text{erf}\zeta \quad \zeta = \frac{Z}{\sqrt{Dt}}$$

$$-(C_A/C_{A0}) + 1 = \text{erf}\zeta$$

$$C_A/C_{A0} = 1 - \text{erf}\zeta$$

$$\text{erf}\zeta = 1 - (0.02/0.05) = 0.6 \rightarrow \zeta \approx 0.6$$

$$\zeta = 0.55 + 0.05 \left(\frac{0.6 - 0.563}{0.604 - 0.563} \right) = 0.595$$

$$0.595 = \frac{10\text{cm}}{\sqrt{2.5 \times 10^{-4} \text{ sec } t}}$$

$$(0.595)^2 = \frac{100}{20 \times 10^{-4} t}$$

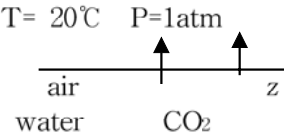
$$t = \frac{100}{(20 \times 10^{-4}) (0.595)^2} = 14100 \text{ sec}$$

or 39 hr
= 1.6 day

More problems to be analyzed by error functions

Carbon dioxide enrichment of a stagnant air mass

A region of air enriched with CO₂ grows upward from the interface as the air mass remains in contact with the water body. Consider the CO₂ distribution in the stagnant air T= 20°C



$$\frac{y_1 - y_{10}}{y_{1\infty} - y_{10}} = \text{erf} \zeta \quad \zeta = \frac{z}{\sqrt{Dt}}$$

$$D = 0.153 \text{ cm}^2 \text{ s} \quad y_{1\infty} = 0.03 \quad y_{10} = 0.0658$$

time	Penetration Distance z (cm)		
	0.1	1	10
10 sec	0.0642	0.0503	0.03
1 min	0.0651	0.0592	0.0307
1 hr	0.0657	0.0649	0.0574

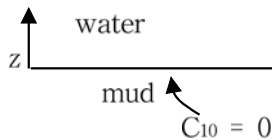
Homework Problem

consider a stagnant layer of O₂-rich water suddenly placed above an anaerobic mud layer capable of consuming O₂ and depleting the water of its O₂ content

Initially water was saturated at 9.17 mg/L

The Diffusivity of O₂ in water is $1.80 \times 10^{-5} \text{ cm}^2 \text{ s}$

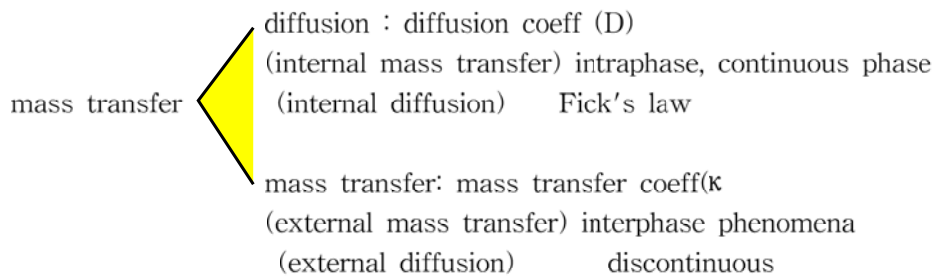
Calculate the concentration of O₂ at given time and distance



time	distance		
	0.001 cm	0.1 cm	10 cm
5 min			
1 hr			
1 day			

Exercises

1. Reverse osmosis takes place in a plate and frame module (between two flat membranes), i.e., the salt is rejected by the membranes and concentrated while water permeate through the membrane. Derive the transport equation that describes the salt concentration, starting from a mass balance on the differential control volume. Assume steady state and no movement in y -direction.
2. A chemical is injected into a fully-developed pipe flow and degraded by a first-order reaction. Starting from mass balance in a differential control volume, derive a differential equation for one-dimensional non-steady transport of the chemical in the pipe in terms of diffusion and bulk velocity. (Boundary conditions not needed)
3. We consider a case where a vessel containing a selective membrane is initially filled to a volume V_0 with solution at a concentration of chemical able to pass through the membrane, c_{r0} . At $t = 0$ fluid begins leaving the vessel at a constant rate Q . The concentration in the effluent (or membrane permeate), c_p , is always less than the concentration in the vessel, c_r , by the constant ratio $p_c = c_p/c_r$ ($p_c < 1$). The effluent is received in another vessel which is well mixed. Derive an expression for the concentration of the chemical in the receiving vessel as a function of time. The receiving vessel is initially empty. If possible, keep the symbols used in Example 1.6, p17, Logan.
4. A stagnant water is saturated with oxygen in air at 25 °C. Suddenly (at time zero) the air is replaced by pure oxygen, and the interface is saturated with the pure oxygen immediately and oxygen diffuses into the water. Calculate the oxygen concentration at 0.5 cm under the interface and the oxygen flux at the interface at 1 hr.



Exercise **Diffusion problems vs. mass transfer problem (D or κ**

(1) catalyst intraparticle phenomena
D

(2) penetration of acid rain into a mable stature
D

(3) O₂ transfer in a activated sludge process
aeration. bubble → bulk O₂ transfer κ

(4) In an activated sludge
nutrient transport into a floc
(substrate) D

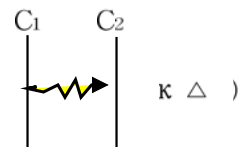
(5) SO₂ scrubbing SO₂ + CaO → CaSO₃
"κ scale up

(6) Cl₂ penetration into bridge

 D

(7) odor control by plastic film
leak or diffusion

mass transfer
flux = κ (concentration difference)



flux variation with position ignored

$$\text{flux} = D \left(\frac{\text{conc difference}}{\text{length}} \right)$$

$$- D \frac{\partial}{\partial x}$$

variation with position predicted

$$= \frac{D}{\ell} (\text{conc difference})$$

"lumped-parameter" vs "distributed parameter"